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Spreading Electron Density Thin: Increasing the Chromophore Size in Polyaromatic Wires Decreases Interchromophoric Electronic Coupling

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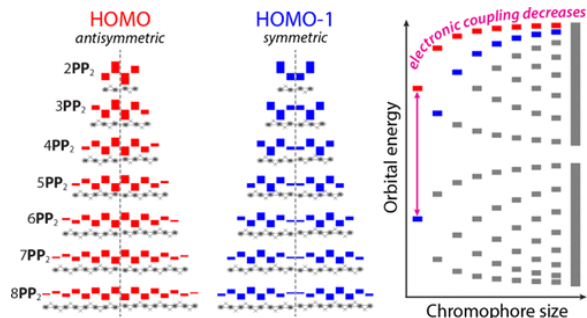
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Abstract



The development of novel polychromophoric materials using extended polycyclic aromatic hydrocarbons as a single large chromophore holds promise for long-range charge-transfer applications in photovoltaic devices and molecular electronics. However, it is not well-understood how the interchromophoric electronic coupling varies with the chromophore size in linearly connected molecular wires. Here, we show with the aid of electrochemistry, electronic spectroscopy, density functional theory calculations, and theoretical modeling that as the number of aromatic moieties in a single chromophore increases, the interchromophoric electronic coupling decreases and may reach negligible values if the chromophore is sufficiently large. The origin of this initially surprising result becomes clear when one considers this problem with the aid of Hückel molecular orbital theory, as at the polymeric limit energies of the molecular orbitals cluster to form bands and thus the energy spacing between orbitals, and thereby the electronic coupling must decrease with the chromophore expansion.

Introduction

Poly-*p*-phenylene-based wires are prototypical systems for charge-transfer studies with potential applications in photovoltaic and molecular electronics devices.^(1–5) Electronic coupling between a pair of phenylenes in a poly-*p*-phenylene-based wire is a crucial parameter that controls its redox and optical properties, as well as the rates of electron transfer in the corresponding donor–wire–acceptor systems.^(6–9) For example, unsubstituted poly-*p*-phenylene wires are characterized by strong interchromophoric electronic coupling because of the favorable nodal arrangement of the highest occupied molecular orbital (HOMO) lobes and relatively small interplanar dihedral angles, which promote effective orbital overlap between adjacent phenylenes.^(10,11) This strong electronic coupling is reflected in the sensitivity of redox/optical properties to the wire length, as can be judged by large slopes in their $1/n$ or $\cos[\pi/(n + 1)]$ dependences,^(10,12) where n is number of chromophoric units in a wire.

In this context, an interesting question concerns how many phenylenes should be included in a single chromophore. For example, a poly-*p*-phenylene-based wire shown in [Figure 1](#) can be either considered as a poly-fluorene (i.e., **PF_n**) or poly-*p*-phenylene (i.e., **PP_n**). Irrespective of how the chromophore is defined, the absorption band of **PF_n** or **PP_n** shifts red with the increasing wire length, indicating that interchromophoric electronic coupling is significant ([Figure 1](#)).⁽¹⁰⁾ However, can the value of the electronic coupling depend on how one defines the chromophore?

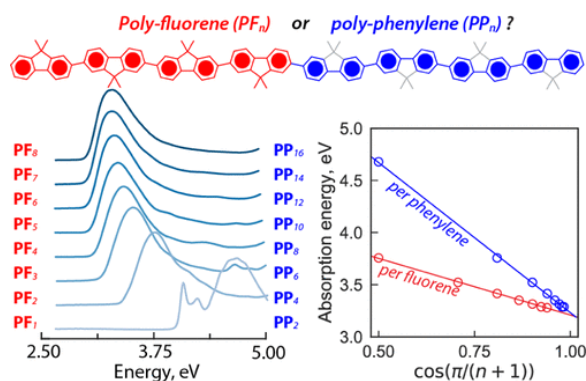


Figure 1. Depending on the choice of the monomeric unit, a molecular wire on the top can be either considered as poly-fluorene (PF_n) or poly-phenylene (PP_n). Left: Absorption spectra of PF_n (or PP_n). Right: Energies of maximum absorption plotted against $\cos[\pi/(n+1)]$, where n is number of phenylenes (blue) or fluorenes (red).

According to the Hückel molecular orbital (MO) theory, the energy of the HOMO to LUMO transition scales linearly with $\cos[\pi/(n+1)]$, and the scaling factor depends on electronic coupling.^(10,13) Remarkably, the slope of the linear ν -vs- $\cos[\pi/(n+1)]$ plot is smaller by nearly a factor of two when n is number of fluorenes as compared to the plot where n is the number of phenylenes (Figure 1), suggesting that the electronic coupling between fluorenes is by a factor of two smaller than that between a pair of phenylenes for the same wire.

As the magnitude of the electronic coupling directly impacts the redox/optical properties of the wire and the extent of hole (i.e., polaron) delocalization in the corresponding cation radicals, a choice of the chromophore size seems to be an important additional parameter. Indeed, a recent study showed on the example of biaryls with varied chromophore size that as the chromophore size increases, the mechanism of hole delocalization changes from static delocalization to dynamic hopping.⁽¹⁴⁾ In order to systematically probe the curious dependence of the chromophore size on the electronic coupling, herein, we designed a set of poly-*p*-phenylene-based mono- and bichromophores (i.e., nPP and nPP_2 , Chart 1) where the size of the chromophore is varied from $n = 2$ to 4 phenylenes. Aided by the electrochemical properties of nPP/nPP_2 and spectroscopic signatures of their cation radicals with the support from density functional theory (DFT) calculations, we conclude that as the size of the chromophore increases, the interchromophoric electronic coupling gradually decreases and may reach nearly nonexistent values at the polymeric limit. We show that the origin of this initially surprising result becomes clear when one considers this problem with the aid of Hückel MO theory, which predicts that at the polymeric limit, the energies of the MOs become nearly degenerate and the energy bands are formed, leading to the inevitability of the decrease of interchromophoric electronic coupling as the chromophore expands.

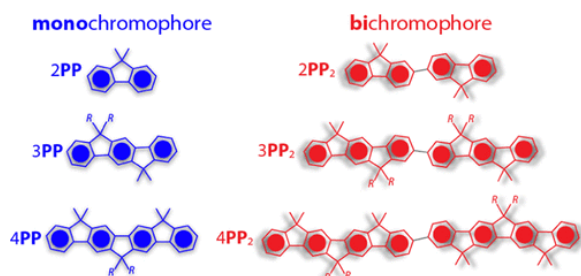


Chart 1. Structures of Monochromophoric nPP and Bichromophoric nPP_2 ($n = 2-4$, $R = \text{Hexyl}$)

Results and Discussion

Synthesis

The preparation of monochromophores $n\text{PP}$ and bichromophores $n\text{PP}_2$ ($n = 3, 4$) was accomplished by adaptation of literature procedures^(15–23) and involved alkylation of 9,9 position of fluorene unit, palladium-catalyzed Suzuki coupling, alkylation of the carbonyl group of the ester to form the tertiary alcohol, and acid-catalyzed Friedel–Crafts intramolecular cyclization. The synthetic procedures are summarized in Scheme S1 in the [Supporting Information](#) along with the full experimental details. All compounds were characterized by $^1\text{H}/^{13}\text{C}$ NMR spectroscopy and matrix-assisted laser desorption ionization (MALDI) mass spectrometry.

Electrochemistry

The electron donor strengths of $n\text{PP}$ and $n\text{PP}_2$ ($n = 2–4$) were evaluated by electrochemical oxidation at a platinum electrode as a 2 mM solution in CH_2Cl_2 containing 0.1 M $n\text{-Bu}_4\text{NPF}_6$ as the supporting electrolyte. The cyclic voltammograms (CVs) of $n\text{PP}$ showed that upon increasing the number of phenylenes, the first oxidation potential decreases from $E_{\text{ox}} = 1.27$ V vs Fc/Fc^+ in 2PP to 0.89 V in 3PP and to 0.68 V in 4PP ([Figure 2](#)). Noteworthy, although 2PP (expectedly) displays an irreversible CV because of the presence of the substitution-labile carbons that render the cation radicals susceptible to dimerization, the CVs of 3PP and 4PP were found to be completely reversible.

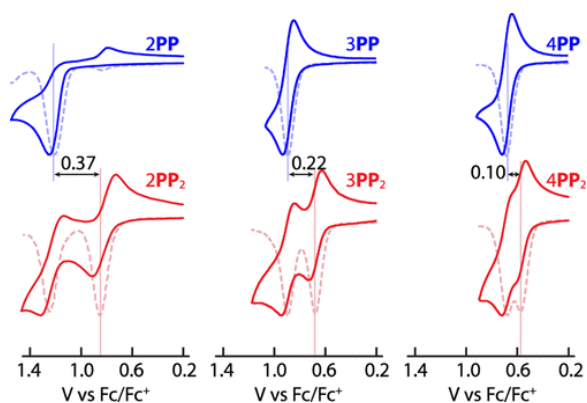


Figure 2. Cyclic (solid lines) and square-wave (dashed lines) voltammograms of 2 mM $n\text{PP}$ (blue) and $n\text{PP}_2$ (red) in CH_2Cl_2 (0.1 M $n\text{-Bu}_4\text{NPF}_6$) at 200 mV s^{-1} and 22 $^\circ\text{C}$. CVs and SWVs of 2PP and 2PP_2 are available from the literature.⁽¹⁰⁾

The reversible CVs of $n\text{PP}_2$ ($n = 2–4$) showed two well-separated oxidation waves that correspond to the formation of the cation radical and dication ([Figure 2](#)). As the number of phenylenes in each chromophore increases, the first oxidation potential of $n\text{PP}_2$ decreases from $E_{\text{ox}} = 0.90$ V vs Fc/Fc^+ in 2PP_2 to 0.67 V in 3PP_2 and to 0.58 V in 4PP_2 . Because the E_{ox} of monochromophoric $n\text{PP}$ decreases *faster* with increasing n than E_{ox} of bichromophoric $n\text{PP}_2$, the amount of hole stabilization, measured as the difference $\Delta E_{\text{ox}} = E_{\text{ox}}[n\text{PP}] - E_{\text{ox}}[n\text{PP}_2]$, decreases from 0.37 to 0.22 to 0.10 V, indicating that the interchromophoric electronic coupling decreases with increasing n ([Table 1](#)). Consistent with this conclusion, we also note that the separation between two oxidation waves in $n\text{PP}_2$ decreases with increasing chromophore size, signifying that the repulsive interactions between the positive charge densities in $n\text{PP}_2^{2+}$ also decrease with increasing chromophore size.^(24–27)

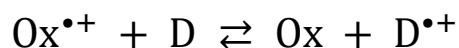
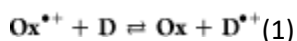
Table 1. Oxidation Potentials (E_{ox} , V vs Fc/Fc^+) of $n\text{PP}$ and $n\text{PP}_2$; Experimental (ΔE_{ox} , V) and Computed (ΔG_{ox} , eV) Using B1LYP40/6-31G(d) + PCM(CH_2Cl_2) Hole Stabilizations; Experimental Wavelengths (λ , nm) of Maximum Absorption of $n\text{PP}^{*+}$ and $n\text{PP}_2^{*+}$ ^a

n	$E_{\text{ox}}[n\text{PP}]$	$E_{\text{ox}}[n\text{PP}_2]$	ΔE_{ox}	ΔG_{ox}	$\lambda[n\text{PP}^{*\cdot}]$	$\lambda[n\text{PP}_2^{*\cdot}]$
2	1.27	0.9	0.37	0.37	690	1240
3	0.89	0.67	0.22	0.19	872	1784
4	0.68	0.58	0.10	0.08	1100	2330

^aValues of $\lambda[2\text{PP}^{*\cdot}]$ and $\lambda[2\text{PP}_2^{*\cdot}]$ are available from literature.⁽¹⁰⁾

Generation of Cation Radicals of $n\text{PP}$ and $n\text{PP}_2$

To further probe the electronic structure of $n\text{PP}_2$, we resorted to electronic spectroscopy of their cation radicals. The cation radicals are expected to display an intervalence transition in the near-IR region, the position and intensity of which are sensitive to the degree of electronic coupling.^(28–30) Cation radicals of $n\text{PP}_2$ and $n\text{PP}$ ($n = 3$ and 4) were generated via quantitative^(31,32) redox titrations using two aromatic oxidants: $[\text{THEO}^{*\cdot}\text{SbCl}_6^-]$ (**THEO** = tetrasubstituted *p*-hydroquinone ether, $E_{\text{red}1} = 0.67$ V vs Fc/Fc⁺, $\lambda_{\text{max}} = 518$ nm, and $\epsilon_{\text{max}} = 7300$ cm⁻¹ M⁻¹)⁽³³⁾ and $[\text{NAP}^{*\cdot}\text{SbCl}_6^-]$ (**NAP** = cycloannulated naphthalene derivative, $E_{\text{red}1} = 0.94$ V vs Fc/Fc⁺, $\lambda_{\text{max}} = 672$ nm, and $\epsilon_{\text{max}} = 9300$ cm⁻¹ M⁻¹).^(34,35) Each redox titration experiment was carried out by an incremental addition of substoichiometric amounts of the electron donor (**D**, i.e., $n\text{PP}_2$ or $n\text{PP}$) to the solution of an oxidant cation radical (**Ox**^{•+}, i.e., **THEO**^{•+} or **NAP**^{•+}). The one-electron oxidation of **D** to **D**^{•+} and reduction of **Ox**^{•+} to **Ox** can be described by an equilibrium shown in [eq 1](#).



The reproducible spectra of $n\text{PP}^{*\cdot}/n\text{PP}_2^{*\cdot}$ were obtained by a numerical deconvolution procedure^(31,32) at each titration point that produced the mole fraction plots of each species involved in the redox reaction. The plots show that, in the case of **3PP**, **4PP**, and **3PP**₂, the addition of a neutral donor leads to exclusive formation of the corresponding cation radical and complete consumption of the oxidant ([Figure 3A–F](#)). In the case of **4PP**₂, the redox titrations involved two successive one-electron oxidations with multiple equilibria ([Figure 3G,H](#)). Upon addition of 0.5 equiv of **4PP**₂, the oxidant is consumed completely and only the dication **4PP**₂²⁺ is formed because of similar values of its first and second oxidation energies ([Table 1](#)). Upon further addition of 0.5 equivalents of neutral **4PP**₂, **4PP**₂²⁺ is completely converted into the **4PP**₂^{•+} via a disproportionation reaction.⁽³¹⁾

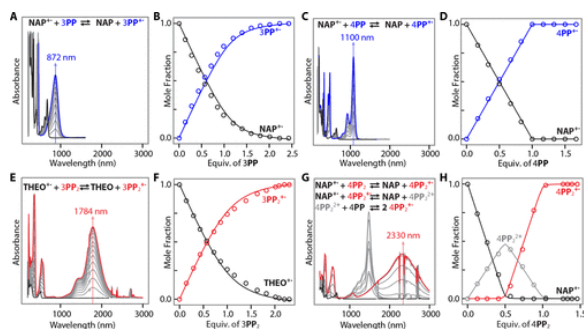


Figure 3. (A) Spectral changes observed upon the reduction of 0.036 mM **NAP**^{•+} in CH₂Cl₂ by incremental addition of 2.2 mM solution of **3PP** in CH₂Cl₂. (B) Mole fraction plot of **NAP**^{•+} (black) and **3PP**^{••} (blue) against the added equivalents of **3PP**. Symbols represent experimental points, whereas the solid lines show best-fit to the experimental points using $\Delta G_1 = -0.08$ V.⁽³¹⁾ (C) Spectral changes observed upon the reduction of 0.021 mM **NAP**^{•+} in CH₂Cl₂ by incremental addition of 2.1 mM solution of **4PP** in CH₂Cl₂. (D) Mole fraction plot of **NAP**^{•+} (black) and **4PP**^{••} (blue) against the added equivalents of **4PP**. Symbols represent experimental points, whereas the solid lines show best-fit to the experimental points using $\Delta G_1 = -1.16$ V. (E) Spectral changes observed upon the reduction of 0.027 mM **THEO**^{•+} in CH₂Cl₂ by incremental addition of 0.77 mM solution of

3PP_2 in CH_2Cl_2 . (F) Mole fraction plot of THEO^{*+} (black) and 3PP_2^{*+} (red) against the added equivalents of 3PP_2 . Symbols represent experimental points, whereas the solid lines show best-fit to the experimental points using $\Delta G_1 = -0.06$ V. (G) Spectral changes observed upon the reduction of 0.018 mM NAP^{*+} in CH_2Cl_2 by incremental addition of 0.45 mM solution of 4PP_2 in CH_2Cl_2 . (H) Mole fraction plot of NAP^{*+} (black), 4PP_2^{2+} (gray), and 4PP_2^{*+} (red) against the added equivalents of 4PP_2 . Symbols represent experimental points, whereas the solid lines show best-fit to the experimental points using $\Delta G_1 = -0.39$ V and $\Delta G_{12} = 0.22$ mV.

Electronic absorption spectra of $n\text{PP}_2^{*+}$ ($n = 3-4$), obtained from the deconvolution, show the presence of a characteristic broad band that is red-shifted (i.e., shifted to longer wavelength) as compared to the corresponding monochromophoric $n\text{PP}^{*+}$ (Figure 3), signifying extensive hole delocalization in $n\text{PP}_2^{*+}$. At the same time, as the number of phenylenes in the chromophore increases, the position of the near-IR band shifts to longer wavelengths, indicating that the electronic coupling is decreasing. To further probe the evolution of electronic coupling with the chromophore size, we resorted to DFT calculations.

Electronic Structure Calculations

Accurate description of the electronic structure of π -conjugated cation radicals is challenging for DFT because of the self-interaction error (SIE), which may lead in unfavorable cases to artificial hole delocalization, artificially low oxidation energies and an incorrect description of the excited states.⁽³⁶⁻³⁹⁾ The SIE can be partially corrected by inclusion of a calibrated amount of the exact Hartree-Fock (HF) exchange term into the hybrid density functional.^(28,40) In our past studies, B1LYP⁽⁴¹⁾ functional with 40% of HF exchange term (i.e., B1LYP-40)⁽⁴²⁾ was introduced, where the amount of HF exchange term was fine-tuned to reproduce experimental oxidation potentials and cation radical excitation energies of poly-*p*-phenylenes with increasing number of *p*-phenylenes. Noteworthy, it was later shown that B1LYP-40/6-31G(d) performs exceptionally well in reproducing the experimental redox/optoelectronic properties of a variety of π -conjugated⁽¹⁰⁻¹²⁾ and π -stacked assemblies⁽⁴²⁻⁴⁴⁾ that were not included in the original training set. Therefore, in this manuscript, we performed DFT calculations of $n\text{PP}/n\text{PP}_2$ and their cation radical using the B1LYP-40/6-31G(d) method and account for the solvent effects using the polarizable continuum model (PCM)⁽⁴⁵⁾ with CH_2Cl_2 parameters.

X-ray crystallography of numerous neutral aromatic hydrocarbons and their cation radicals has established that oxidation induces significant structural reorganization in the form of elongations and contractions of the C-C bonds, as well as in the decrease of dihedral angles between adjacent aromatic moieties.^(46,47) Indeed, calculations showed that upon oxidation, the molecular structure of $n\text{PP}_2$ undergoes a quinoidal distortion, as exemplified for 2PP_2 in Figure 4A below.

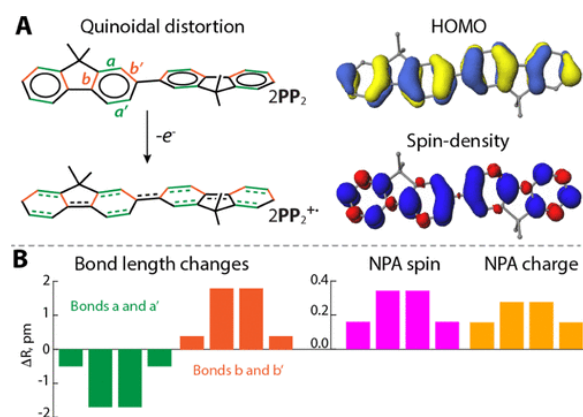


Figure 4. (A) Schematic representation of the quinoidal distortion in $2\text{PP}_2 \rightarrow 2\text{PP}_2^{+\bullet}$ transformation, HOMO of 2PP_2 , and spin density of $2\text{PP}_2^{+\bullet}$ calculated using B1LYP-40/6-31G(d) + PCM(CH_2Cl_2). (B) Per-phenylene bar-plot representation of the distributions of the bond length changes, NPA spin, and charge distributions in $2\text{PP}_2^{+\bullet}$.

Importantly, oxidation-induced bond length changes track in accordance with the disposition of the lobes of HOMO, that is, bonds that correspond to the antibonding lobes undergo contraction and bonds that correspond to the bonding lobes undergo elongations, for example, the average of bonds a and a' contracts by 0.5–1.7 pm, whereas the average of bonds b and b' elongates by 0.4–1.8 pm (Figure 4). The natural population analysis (NPA) of the electron density in $2\text{PP}_2^{+\bullet}$ further showed that the spin and charge distributions parallel the distribution of the bond-length changes, indicating polaron formation^(48–50) (Figure 4B).

A simple yet intuitive model to describe the polaron (i.e., hole) delocalization in various mixed-valence compounds is the two-state Marcus–Hush theory, which predicts that the extent of polaron delocalization depends on the interplay between electronic coupling (H_{ab}) and reorganization energy (λ).^(28,29,51,52) In the large electronic coupling limit ($2H_{ab} \geq \lambda$), the polaron is evenly delocalized between two sites (i.e., class III or “static” delocalization), whereas in the limit of small electronic coupling ($2H_{ab} < \lambda$), the polaron is partially delocalized with the distribution maximum centered on one of the sites (i.e., class II or “dynamic” hopping). When electronic coupling is nonexistent ($H_{ab} = 0$, class I), the polaron is fully localized on one site.

Partitioning the spin, charge, and structural reorganization distributions in $2\text{PP}_2^{+\bullet}$ between a pair of fluorenes rather than four phenylenes suggests that the polaron is evenly delocalized between the fluorenes, and hence the mechanism of polaron delocalization is “static” delocalization according to the two-state representation (Figure 5).

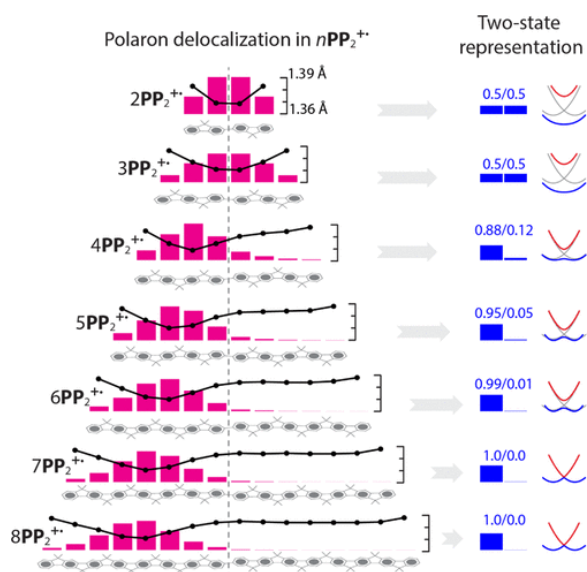


Figure 5. Left: Per-phenylene bar-plot representation of the distributions of the NPA spin (magenta) superimposed with the plot of C–C bond length (average of a and a' , Figure 4) in $n\text{PP}_2^{+\bullet}$. Right: Two-state representation of polaron delocalization in $n\text{PP}_2^{+\bullet}$.

Similarly, in $3\text{PP}_2^{+\bullet}$, the spin/charge/bond length changes are delocalized along the entire bichromophore, and thus the mechanism of polaron delocalization is also static delocalization. However, further increase of the chromophore size shifts the polaron distribution toward one side of $n\text{PP}_2^{+\bullet}$ ($n > 3$), indicating that the mechanism of polaron delocalization evolves into dynamic hopping, that is, class II (Figure 5). In the case of even

longer nPP_2^{*+} ($n > 5$), the polaron is fully localized on a single chromophore, that is, this bichromophore belongs to class I. Thus, depending on the chromophore size, the mechanism of hole delocalization in a two-state representation may switch from static delocalization to dynamic hopping to a complete localization, akin to the switchover observed for biaryls with varied interplanar angle.⁽⁵²⁾ For the latter, an increase in the interplanar angle led to the decrease in the overlap between p-orbitals of carbon atoms at the biaryl linkage, leading to the reduced electronic coupling. As the extent of hole delocalization depends on the interplay between electronic coupling (H_{ab}) and reorganization energy (λ), the switchover in the delocalization mechanism in nPP_2^{*+} must arise because of the decreasing H_{ab} and/or increasing λ .⁽⁵¹⁾

To confirm that the switchover in the delocalization mechanism occurs because of the decreasing electronic coupling, we first compare computed oxidation free energies (G_{ox}) of nPP and nPP_2 ($n = 2-8$). Calculations showed that as the number of phenylenes in the chromophore increases, oxidation energies of both nPP and nPP_2 decrease (Figure 6A). Furthermore, hole stabilization, measured as the difference $\Delta G_{ox} = G_{ox}[nPP] - G_{ox}[nPP_2]$, is the largest for $n = 2$ and decreases with increasing n . Comparing the computed ΔG_{ox} with the available experimental values of ΔE_{ox} for $n = 2-4$ (Table 1), it is clear that the B1LYP-40 method shows a remarkable performance in reproducing hole stabilization for these systems, as the error in the computed ΔG_{ox} values is less than 0.03 eV. This suggests that calculations can provide reliable information on the extent of hole stabilization for the longer bichromophores, for which experimental data are unavailable. The calculations show that for longer nPP_2 the increase in chromophore size reduces the hole stabilization to nearly nonexistent values, that is, $\Delta G_{ox} < 0.01$ for $n > 5$, which is consistent with complete localization of the polaron on a single chromophore.

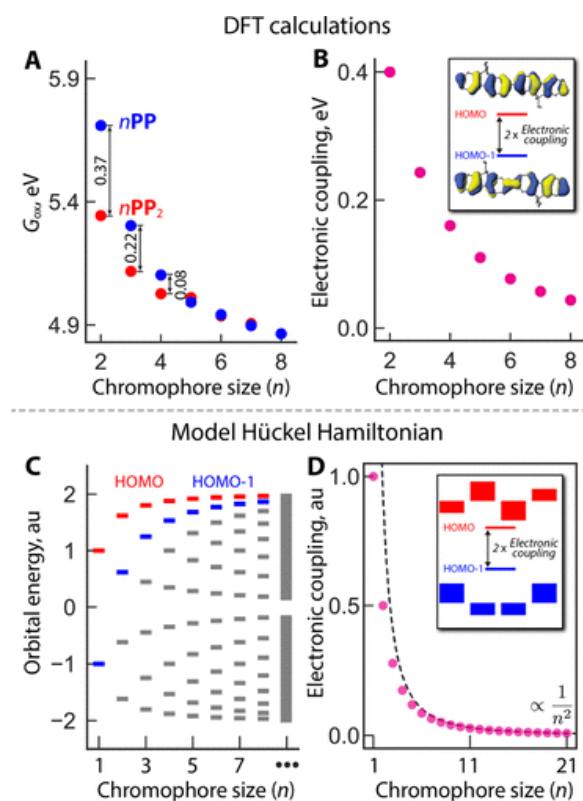


Figure 6. (A) Oxidation energies (G_{ox}) of nPP (blue) and nPP_2 (red) against the chromophore size measured in the number of phenylene units (n) computed using B1LYP-40/6-31G(d) + PCM(CH_2Cl_2). (B) Interchromophoric electronic coupling of nPP_2 measured as a half of the HOMO/HOMO – 1 energy gap against chromophore size (n). Inset shows HOMO and HOMO – 1 of $2PP_2$. (C) Orbital energies calculated for a model Hückel Hamiltonian matrix (eq 2) against number of phenylene units in a chromophore. (D) Electronic coupling measured as a HOMO/HOMO – 1 energy gap from a model Hückel Hamiltonian matrix against number of phenylene units in a

chromophore. Dashed line corresponds to the $3\pi^2/8n^2$ dependence derived in the limit of $n \rightarrow \infty$; see the [Supporting Information](#) for details. Inset shows per-phenylene bar-plot representation of HOMO and HOMO – 1 wavefunctions of 2PP_2 obtained from the Hückel Hamiltonian model.

MO Theory of $n\text{PP}_2$

Following Hückel's approach of representing MOs as a linear combination of atomic orbitals, the two upper occupied frontier MOs of $n\text{PP}_2$ obtained from DFT calculations can be approximated as symmetric and antisymmetric linear combinations of the HOMO of a single chromophore (Figure S5 in the [Supporting Information](#)). Then, the energy difference between HOMO and HOMO – 1 equals twice the electronic coupling and thus can be quickly determined from the DFT calculations of neutral $n\text{PP}_2$. Following this approach, electronic couplings were computed for each $n\text{PP}_2$, which showed a decrease with an increasing chromophore size ([Figure 6B](#)).

To provide an intuitive rationale for this observation, we resorted to a theoretical model. Following Hückel theory, each bichromophore $n\text{PP}_2$ can be represented as a set of $2n$ electronically coupled phenylenes with the Hamiltonian matrix **H**:

$$\mathbf{H} = \begin{bmatrix} \alpha & \beta & \cdots & 0 & 0 \\ \beta & \alpha & \cdots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \cdots & \alpha & \beta \\ 0 & 0 & \cdots & \beta & \alpha \end{bmatrix} \quad (2)$$

$$H = \begin{bmatrix} \alpha & \beta & \cdots & 0 & 0 \\ \beta & \alpha & \cdots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \cdots & \alpha & \beta \\ 0 & 0 & \cdots & \beta & \alpha \end{bmatrix}$$

where **H** is the $2n \times 2n$ tight-binding matrix, α is the orbital energy of a single phenylene, and β is the electronic coupling between a pair of adjacent phenylenes. Note that analytical expressions of eigenvalues (i.e., MO energies) and eigenvectors (i.e., MO wavefunctions) of the Hückel Hamiltonian matrix are known from the original works.^(53,54) Here, for the sake of simplicity, we set $\alpha = 0$ and $\beta = -1$.

Diagonalization of the Hamiltonian **H** for varied n produced eigenvalues and eigenvectors, among which two largest eigenvalues correspond to the energies of HOMO and HOMO – 1 ([Figure 6C](#)). The corresponding eigenvectors can be viewed as symmetric and antisymmetric linear combinations of the HOMO of a single chromophore ([Figure 7](#)). Thus, based on the simple Hückel Hamiltonian approach, the energy gap between HOMO and HOMO – 1 decreases with increasing number of phenylenes in a chromophore, leading to nearly isoenergetic values and thus negligible electronic coupling at large n ([Figure 6D](#)). This is fully consistent with the results obtained from DFT calculations ([Figure 6B](#)). In contrast to DFT calculations, solution to the Hückel Hamiltonian matrix is available in an analytic form^(53,54) and thus allows a derivation of the analytic formula for the electronic coupling, which in the limit of $n \rightarrow \infty$ follows a simple $3\pi^2/8n^2$ dependence (see [Figure 6D](#) and the [Supporting Information](#) for details).

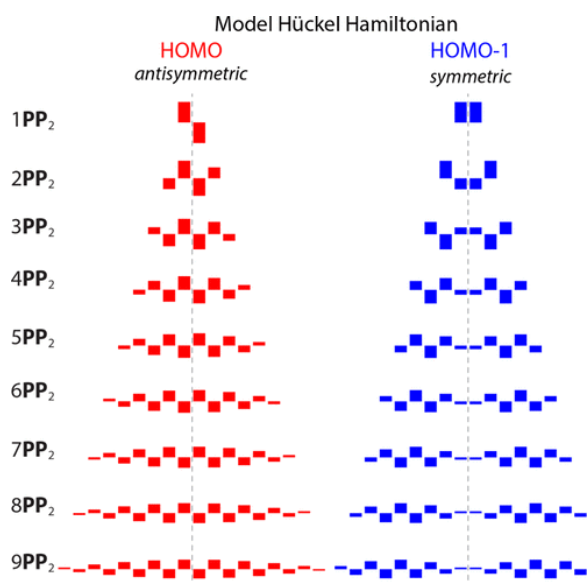


Figure 7. Bar-plot representations of the HOMO and HOMO – 1 wavefunctions of nPP_2 obtained from the diagonalization of the Hückel Hamiltonian model.

We note that although in a discrete molecule, the MO energies are clearly distinguished; at the polymeric limit, the energy *bands* are formed that can be described by a continuous density of states ([Figure 6B](#)).^(55,56) It is thus inevitable that as the oligomer size increases, the spacing between its MO energies (HOMO and HOMO – 1 included) will decrease and reach near-degenerate values at the polymeric limit, which in the case of a bichromophoric oligomer can be interpreted as decreasing the electronic coupling between chromophores.

Following a recent study⁽¹⁴⁾ on the mechanism of hole delocalization in biaryls with varied donor strength of the substituent and varied size of the chromophore, an alternative explanation of the decreasing electronic coupling in nPP_2 is warranted. This study showed that upon increase in the chromophore size, a total of two electrons per orbital spreads over a larger area, thus decreasing the amount of the electron density at the coupling-mediating carbons at the biaryl linkage and leading to the diminished orbital overlap and thereby diminished electronic coupling. Likewise, as the number of phenylenes in each chromophore of nPP_2 increases, the electron density of HOMO spreads over a larger number of carbons, decreasing the electron density at the coupling-mediating carbons and thereby decreasing interchromophoric electronic coupling ([Figure 7](#)).

Conclusions

Motivated by the question whether a set of linearly connected fluorenes is better termed a poly-fluorene or poly-phenylene wire ([Figure 1](#)), we performed a combined experimental/theoretical study on a series of phenylene-based bichromophores nPP_2 of increasing size. We found experimentally that as the number of phenylenes in a chromophore increases from $n = 2$ to 4, the amount of hole stabilization measured electrochemically decreases and the absorption band in the electronic spectra of nPP_2^{*+} shifts to longer wavelength, indicating that the interchromophoric electronic coupling decreases. Aided by the benchmarked DFT calculations of a series of long nPP/nPP_2 ($n = 2-8$), we have shown that the electronic coupling decreases to nearly nonexistent values when $n > 5$, leading to a complete localization of a polaron (i.e., hole) on a single chromophore in the corresponding nPP_2^{*+} . Finally, using the Hückel Hamiltonian model, we showed that the decreasing electronic coupling is a consequence of the clustering of MO energy level as the length of the oligomer increases. This obvious, yet fundamental, understanding that the interchromophoric electronic coupling decreases with increasing chromophore size is crucial during the rational design of novel molecular

wires with large chromophores, where long-range charge transfer can be mediated by incoherent hopping mechanisms.

Supporting Information

The Supporting Information is available free of charge on the [ACS Publications website](https://pubs.acs.org) at DOI: [10.1021/acs.jpcc.8b05299](https://doi.org/10.1021/acs.jpcc.8b05299).

- Experimental and computational details and Hückel theory modeling ([PDF](#))
- Equilibrium geometries of $n\text{PP}$ and $n\text{PP}_2$ ($n = 2-8$) optimized at neutral and cation radical states ([TXT](#))
- pdf
 - [jp8b05299_si_001.pdf \(14.9 MB\)](#)
- text
 - [jp8b05299_si_002.txt \(151.16 kb\)](#)

The authors declare no competing financial interest.

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